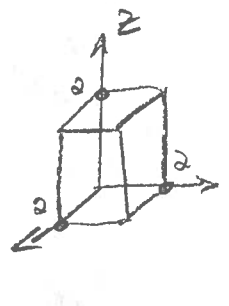


Ryan Chilton
ECE 311 - HW #6

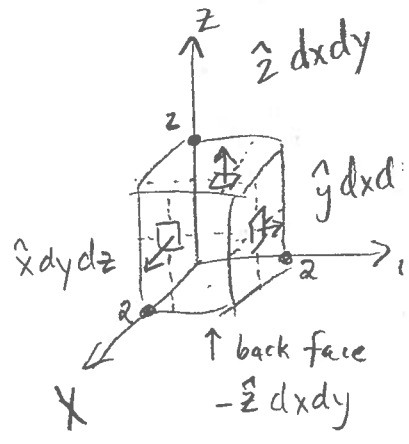
(4.20) (a) $\rho_v = \nabla \cdot \vec{D} = \frac{\partial}{\partial x} D_x + \frac{\partial}{\partial y} D_y + \frac{\partial}{\partial z} D_z$
 $= \frac{\partial}{\partial x} (2x+2y) + \frac{\partial}{\partial y} (3x-2y)$
 $= 2 - 2 = 0 \text{ C/m}^3$

(b) $Q_{total} = \iiint \rho_v dv = \int_0^2 \int_0^2 \int_0^2 0 dx dy dz$
 $= 0 \text{ C}$



(c) $Q_{total} = \oiint_S \vec{D} \cdot d\vec{s}$

Break this into 6 surface integrals:



$Q = \int_0^2 \int_0^2 \vec{D} \cdot \hat{z} dx dy \Big|_{z=2m} +$
 $-\int_0^2 \int_0^2 \vec{D} \cdot \hat{z} dx dy \Big|_{z=0m} +$

$\int_0^2 \int_0^2 \vec{D} \cdot \hat{y} dx dz \Big|_{y=2m} - \int_0^2 \int_0^2 \vec{D} \cdot \hat{y} dx dz \Big|_{y=0m}$

$+ \int_0^2 \int_0^2 \vec{D} \cdot \hat{x} dy dz \Big|_{x=2m} - \int_0^2 \int_0^2 \vec{D} \cdot \hat{x} dy dz \Big|_{x=0m}$

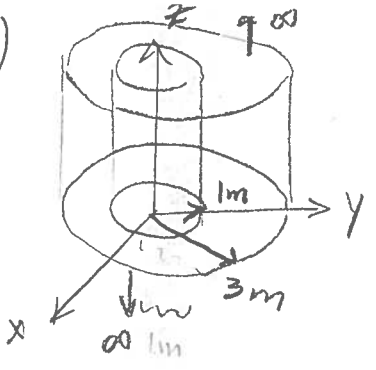
Here, \vec{D} has no \hat{z} component, so the first 2 integrals vanish!

$Q = \int_0^2 \int_0^2 (3x-2y) dx dz \Big|_{y=2m} - \int_0^2 \int_0^2 (3x-2y) dx dz \Big|_{y=0}$
 $+ \int_0^2 \int_0^2 (2x+2y) dy dz \Big|_{x=2m} - \int_0^2 \int_0^2 (2x+2y) dy dz \Big|_{x=0}$

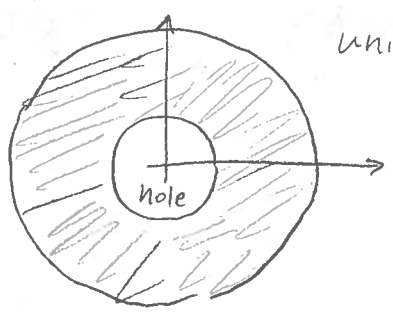
$Q = \int_0^2 \int_0^2 (3x-4) dx dz - \int_0^2 \int_0^2 3x dx dz$
 $+ \int_0^2 \int_0^2 (4+2y) dy dz - \int_0^2 \int_0^2 2y dy dz \quad = -4-12+24-8=0!$

$= 2 \cdot \left(\frac{3}{2}x^2 - 4x \right) \Big|_0^2 - 2 \cdot \left(\frac{3}{2}x^2 \right) \Big|_0^2 + 2(4y+y^2) \Big|_0^2 - 2(y^2) \Big|_0^2$
 $= -4 - 12 + 24 - 8 = 0$

(4.25)



Top view:



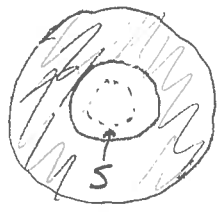
uniform charge density ρ_v

By symmetry considerations, \vec{D} only has an \hat{r} component, and is invariant under z translation and xy rotation (that is, no dependence on ϕ or z). We write:

$$\vec{D} = \hat{r} D_r(r)$$

There's 3 cases to consider as Gaussian Surfaces, S ...

FIRST
 $r < 1m$



$$\oiint \vec{D} \cdot d\vec{s} = Q_{\text{inside}} \quad \left. \begin{array}{l} \text{only cylinder wall} \\ \text{contributes} \end{array} \right\}$$

$$\int_0^h \int_0^{2\pi} \hat{r} D_r \cdot \underbrace{\hat{r}(r d\phi dz)}_{d\vec{s}} = 0 \quad (\text{no charge!})$$

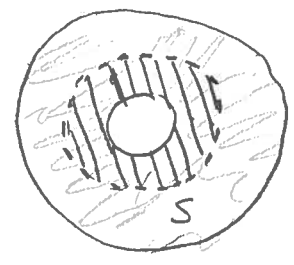
$$\int_0^h \int_0^{2\pi} D_r \cdot r d\phi dz = 0$$

$$D_r \cdot 2\pi r h = 0$$

$$D_r = 0 \quad r < 1m$$

SECOND:

$1m < r < 3m$



$$\oiint \vec{D} \cdot d\vec{s} = Q_{\text{inside}}$$

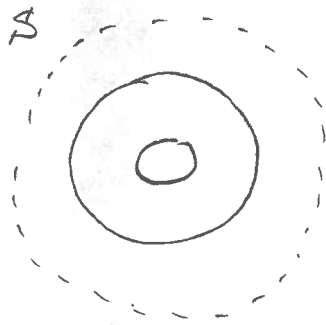
$$\int_0^h \int_0^{2\pi} D_r \cdot r d\phi dz = \rho_v \left[\underbrace{\pi r^2 h}_{\text{uniform charge density}} - \underbrace{\pi (1m)^2 h}_{\text{volume of hatched shell}} \right]$$

$$D_r \cdot 2\pi r h = \rho_v \pi h (r^2 - 1m^2)$$

$$D_r = \frac{\rho_v \cdot (r^2 - 1m^2)}{2r} \quad 1m < r < 3m$$

C/mz

THIRD



$r > 3m$

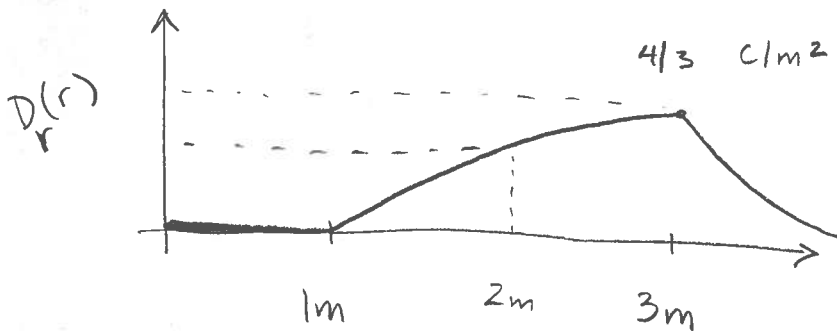
$$\oint \vec{D} \cdot d\vec{s} = Q_{\text{inside}}$$

$$\int_0^h \int_0^{2\pi} D_r \cdot r d\phi dz = \rho_{v0} [\pi(3m)^2 h - \pi(1m)^2 h]$$

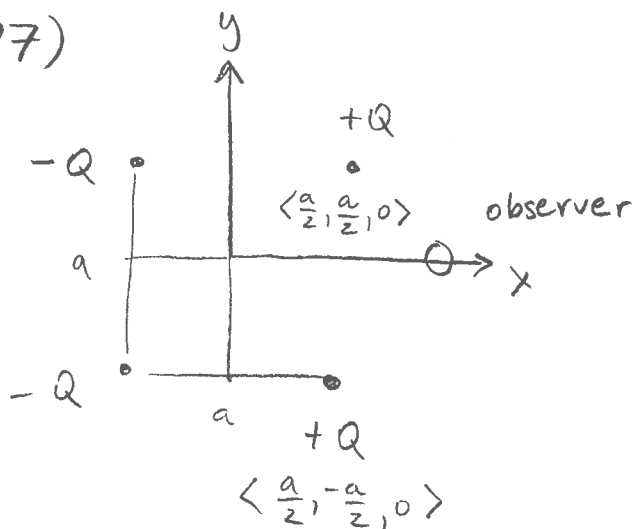
$$2\pi r h \cdot D_r = (\rho_{v0} \cdot 8\pi h) \text{ Coulombs}$$

$$D_r = \frac{\rho_{v0} \cdot 4}{r} \text{ C/m}^2 \quad r > 3m$$

Here's a sketch:



(4.27)



Define a point on the x-axis:

$$\vec{R} = \langle x, 0, 0 \rangle$$

Sum over discrete charges, adding up voltage contributions:

$$V_{\text{total}} = \sum_{i=1}^4 \frac{1}{4\pi\epsilon_0} \frac{q_i}{|\vec{R} - \vec{R}_i|}$$

(a)

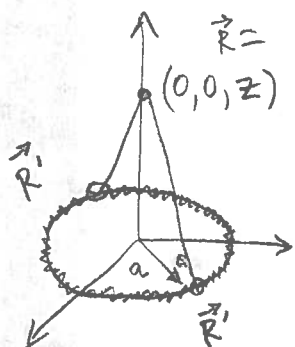
$$V_{\text{total}}(x) = \frac{1}{4\pi\epsilon_0} \left(\frac{Q}{\sqrt{(x - \frac{a}{2})^2 + (\frac{a}{2})^2}} + \frac{Q}{\sqrt{(x - \frac{a}{2})^2 + (\frac{a}{2})^2}} + \frac{-Q}{\sqrt{(x + \frac{a}{2})^2 + (\frac{a}{2})^2}} + \frac{-Q}{\sqrt{(x + \frac{a}{2})^2 + (\frac{a}{2})^2}} \right)$$

(b) Evaluate at $x = \frac{a}{2}$

$$V_{\text{total}}\left(\frac{a}{2}\right) = \frac{1}{4\pi\epsilon_0} \left(\frac{Q}{(a/2)} + \frac{Q}{(a/2)} + \frac{-Q}{\sqrt{a^2 + a^2/4}} + \frac{-Q}{\sqrt{a^2 + a^2/4}} \right)$$

$$= \frac{1}{4\pi\epsilon_0} Q \left(\frac{2}{a} + \frac{2}{a} + \frac{-2}{a\sqrt{5}} + \frac{-2}{a\sqrt{5}} \right) = \frac{Q}{\pi\epsilon_0} \left(1 - \frac{1}{\sqrt{5}} \right)$$

(4.29) Circular ring of charge with radius a , observed at a point $(0, 0, z)$:



Note as we walk around the ring, the distance from the differential piece of charge to the observer doesn't change!
(it's always $\sqrt{z^2 + a^2}$)

Use: $V = \frac{1}{4\pi\epsilon_0} \int_L \frac{\rho \ell}{|\vec{R} - \vec{R}'|} dL'$ over the circular path L' ...

$$V = \frac{1}{4\pi\epsilon_0} \int_{\phi=0}^{\phi=2\pi} \frac{\rho \ell}{\sqrt{z^2 + a^2}} \cdot r d\phi$$

walking around the circle

$\rho \ell / |\vec{R} - \vec{R}'|$

Jacobian

$$V = \frac{1}{4\pi\epsilon_0} \cdot \int_{\phi=0}^{\phi=2\pi} \frac{\rho \ell}{\sqrt{z^2 + a^2}} \cdot a d\phi$$

integration path is constant r ($r = a$)

$$V = \frac{1}{4\pi\epsilon_0} \cdot \frac{\rho \ell \cdot a}{\sqrt{z^2 + a^2}} \left(\int_0^{2\pi} d\phi \right) = \frac{1}{2\epsilon_0} \cdot \frac{\rho \ell \cdot a}{\sqrt{z^2 + a^2}}$$

(b) Compute the -gradient of V to find \vec{E} . Note that \vec{V} is only a function of z ...

$$\vec{E} = -\nabla V = \hat{z} \cdot -\frac{\partial}{\partial z} [V(z)]$$

$$= -\hat{z} \cdot \frac{\rho \ell \cdot a}{2 \cdot \epsilon_0} \cdot \frac{\partial}{\partial z} \left[\frac{1}{\sqrt{z^2 + a^2}} \right]$$

$$= +\hat{z} \frac{\rho \ell \cdot a}{2\epsilon_0} \cdot \frac{-z}{(z^2 + a^2)^{3/2}}$$

$$\frac{\partial}{\partial z} \left[\frac{1}{\sqrt{z^2 + a^2}} \right] = \frac{\partial}{\partial z} (z^2 + a^2)^{-1/2}$$

$$= -\frac{1}{2} (z^2 + a^2)^{-3/2}$$

$$\cdot \frac{\partial}{\partial z} (z^2 + a^2)$$

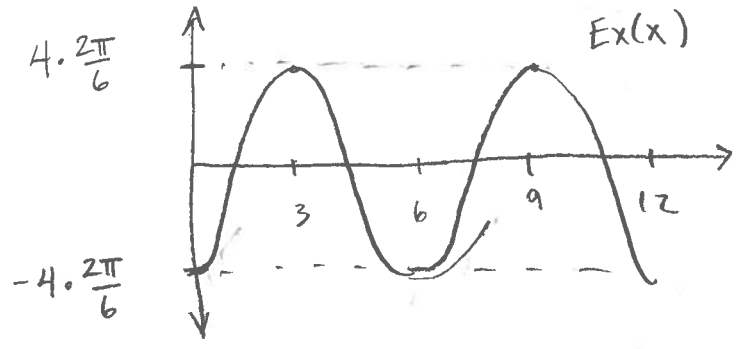
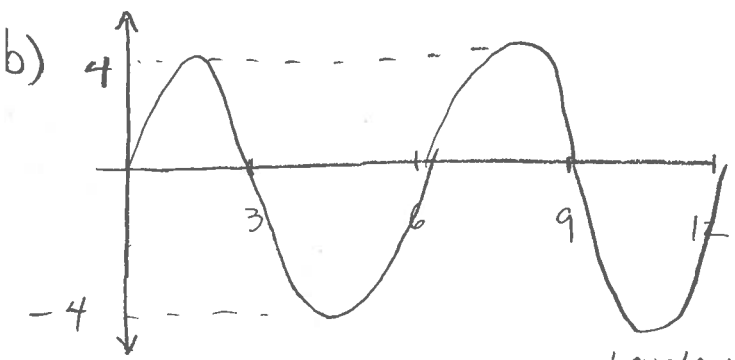
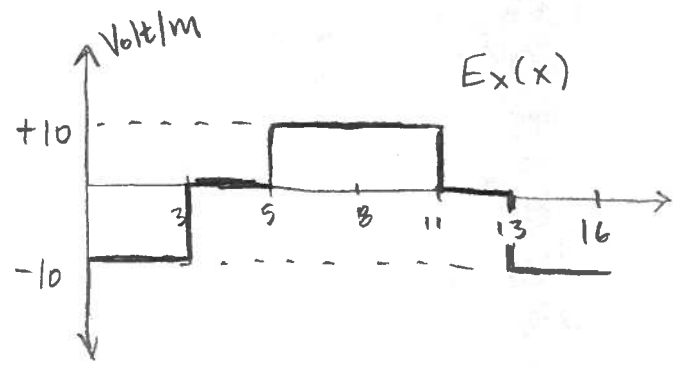
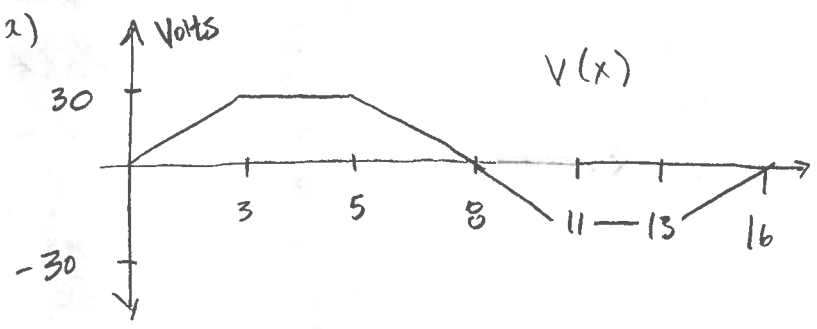
$$= -\frac{1}{2} (z^2 + a^2)^{-3/2} \cdot (2z)$$

$$= -z (z^2 + a^2)^{-3/2}$$

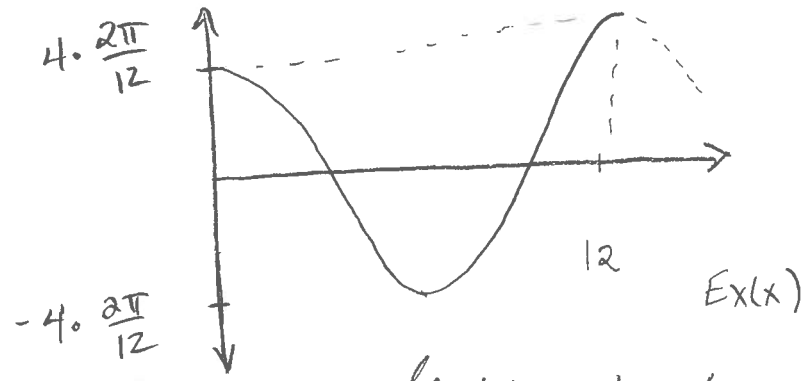
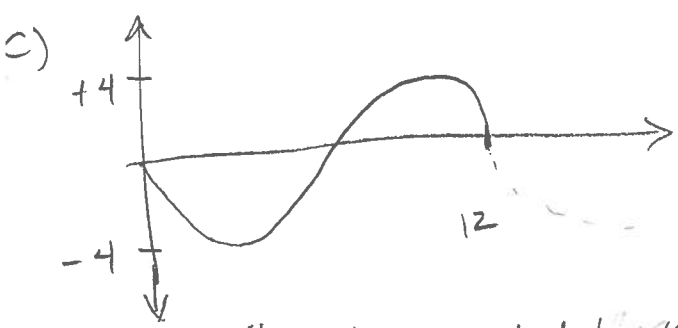
CHAIN RULE + POWER RULE

(4.33) Given the potential graph, plot the electric field.

$$\vec{E} = -\nabla[V(x)] = \hat{x} \cdot \underbrace{-\frac{\partial}{\partial x} V(x)}_{\text{"E}_x\text{"}}$$



period = 6m, freq = 1 cycle/6m = 2π rad/6m
 sinusoid - differentiates to cosine,
 Scaled by frequency



Now the period is 12m,
 and we flipped the sign...

lousy artwork...